

ANALYTICAL STUDIES OF MHD OUTFLOWS

Self-similar solutions

Outflows from stellar bodies

From Solar Wind

To Stellar winds (O M stars)

Compact objects (pulsars, ...)

AGN

GRB

From 1-D Hydro steady model (**Parker 1958**)

To 3-D MHD relativistic time dep. codes

Collimated flows: **Steady** and **axisymmetric**
MHD approach

Steadiness and axial symmetry ($\partial/\partial t = \partial/\partial \phi = 0$)

$$[\mathbf{B}, \mathbf{V}] = [\mathbf{B}_p, \mathbf{V}_p] + [B_\phi, V_\phi] \mathbf{e}_\phi$$

- \mathbf{B}_p and $\rho \mathbf{V}_p$ expressed through the functions $\Psi(r, \theta)$ and $A(r, \theta)$ ($\Psi_A = d\Psi/dA$):

$$\mathbf{B}_p = \nabla \times \left[\frac{A}{r \sin \theta} \mathbf{e}_\phi \right], \quad \rho \mathbf{V}_p = \nabla \times \left[\frac{\Psi}{r \sin \theta} \mathbf{e}_\phi \right]$$

$$\mathbf{V}_p = \frac{\Psi_A}{4\pi\rho} \mathbf{B}_p, \quad M^2 = \frac{V_p^2}{V_{a,p}^2} \equiv \frac{\Psi_A^2}{4\pi\rho}$$

- B_ϕ, V_ϕ expressed in integral form:

$$B_\phi \propto \frac{1 - [\Omega r^2 \sin^2 \theta]/L}{1 - M^2}$$

$$V_\phi \propto \frac{[\Omega r^2 \sin^2 \theta]/L - M^2}{1 - M^2}$$

- $L(A) \propto$ total angular momentum
- $\Omega(A) \propto$ angular velocity of the fieldlines
- $r^2 \sin^2 \theta|_{M=1} = L/\Omega$

Bernoulli + Transfield \rightarrow

$$(1 - M^2) \left(a_1 \frac{\partial^2 A}{\partial r^2} + \frac{2a_2}{r} \frac{\partial^2 A}{\partial r \partial \theta} + \frac{a_3}{r^2 \sin^2 \theta} \frac{\partial^2 A}{\partial \theta^2} + a_4 \right) + a_5 = 0$$

where a_i are complicated functions
The nature of this Eq. ruled by:

$$D = a_2^2 - a_1 a_3 = \\ = (V_p^2 - V_f^2)(V_p^2 - V_s^2)(V_p^2 - V_c^2)(V_a^2 + C_s^2)$$

- $D < 0$: *elliptic*, Dirichlet B.C. on a *closed* surface
- $D > 0$: *hyperbolic*, Cauchy B.C. on an *open* surface
- Our Eq. is of *mixed* type: in different regions we can have $D > 0$ or $D < 0$
- Singularities for $M^2 = 1$ and $V_p = V_s, V_f$
- *Relation between the singularities and the hyperbolic/elliptic regions ?*

Singularities and hyperbolic/elliptic regimes

Slope of the characteristics in the hyp. regions

$$r \frac{d\theta}{dr} = \frac{1}{a_1} (a_2 \pm \sqrt{D})$$

For $a_1 = 0$ one of the characteristics direction becomes parallel to \mathbf{e}_θ :

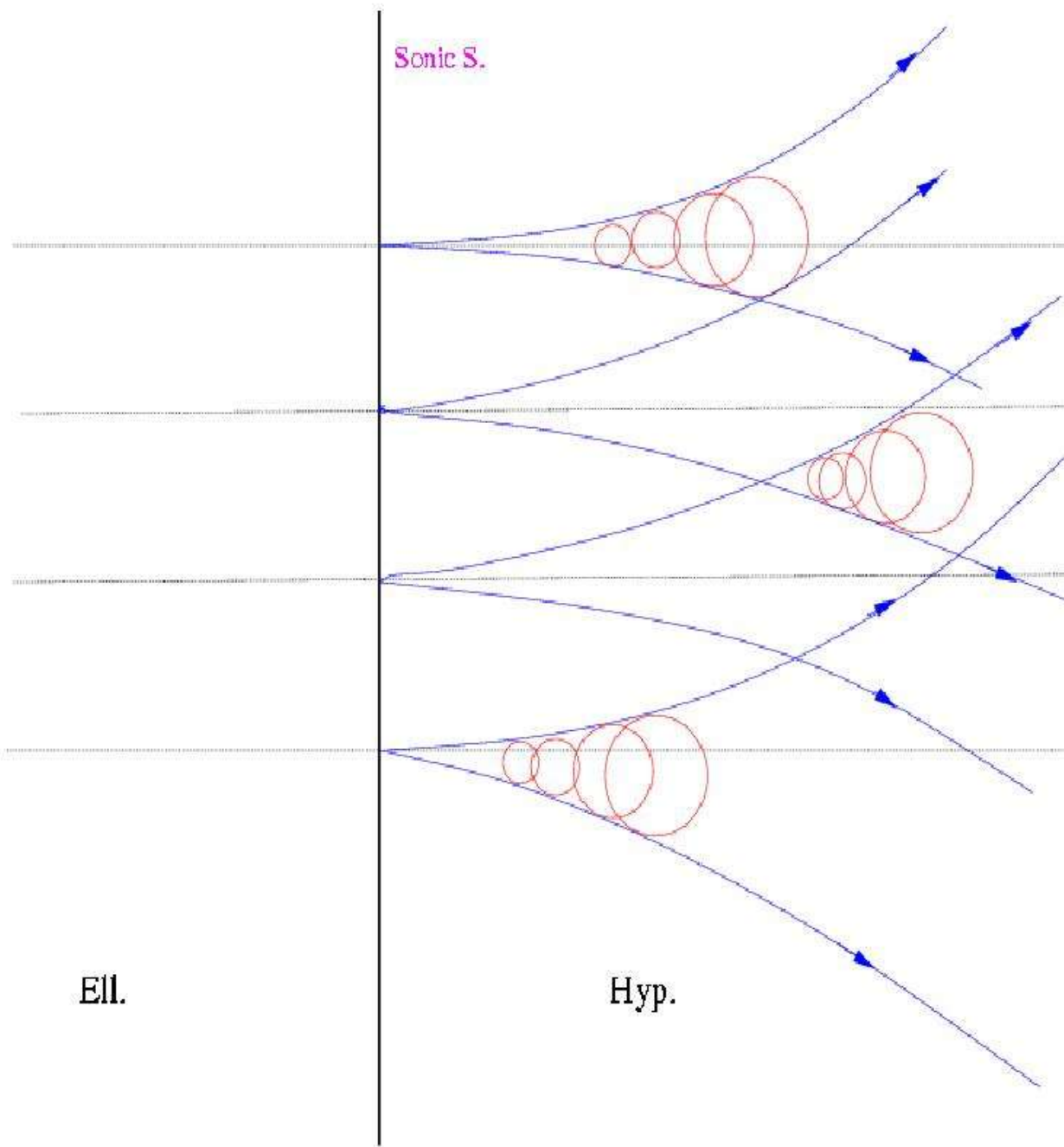
Limiting Characteristics or *Separatrix*

- If the surface of the *L.C.* is not \perp to \mathbf{V}_p it *does not coincide* with the positions where $D = 0$
- The *L.C.* defines the regions *causally connected* inside the outflow

★ Classical hydrodynamical problem in 2-D supersonic flows (Landau & Lifshitz 1959)

To construct a correct solution we need to know the *L.C.*, but *this would require to know this solution !*

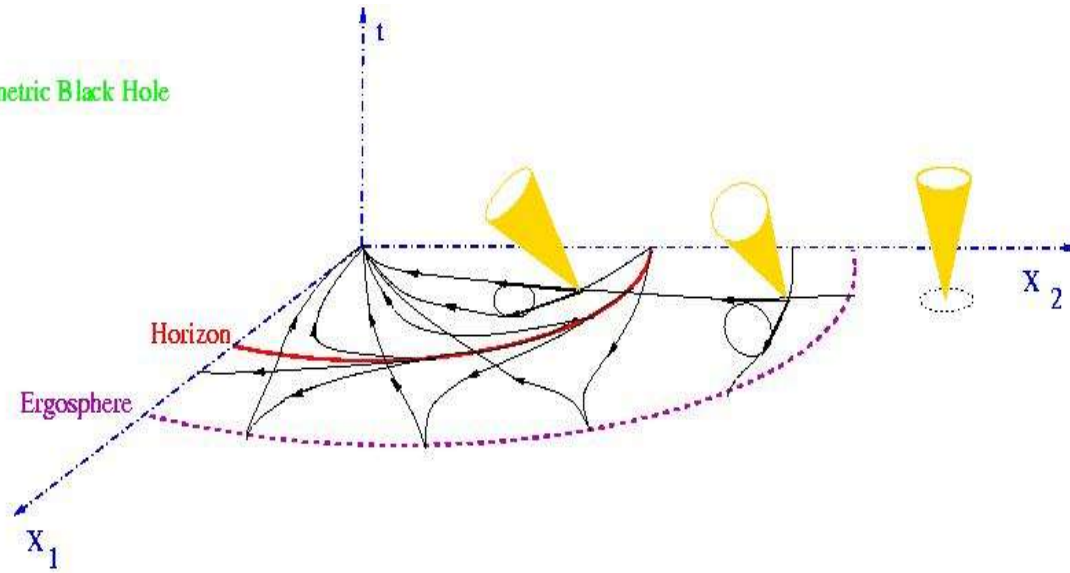
Sonic S.



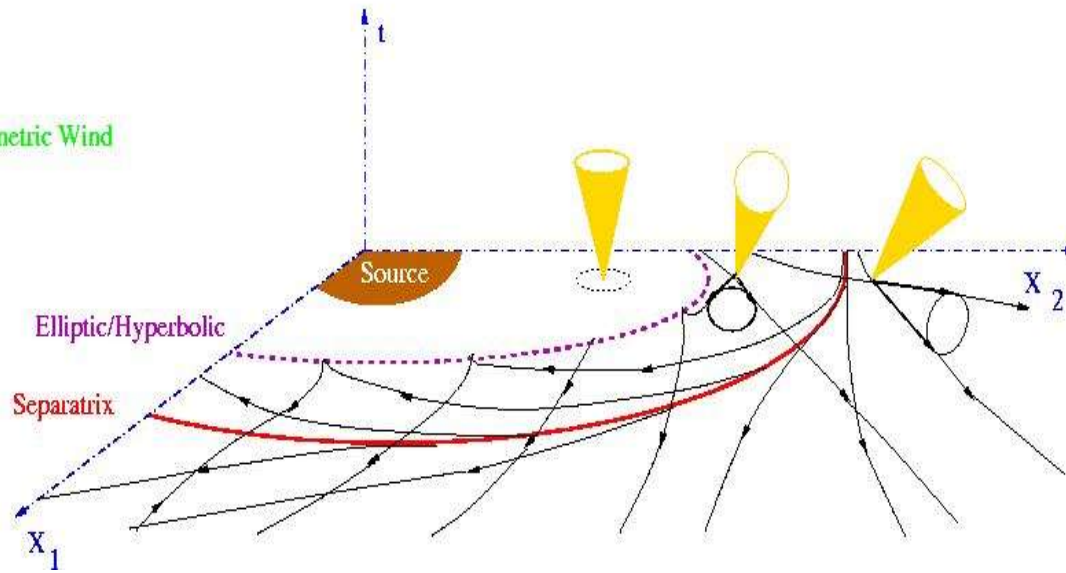
Ell.

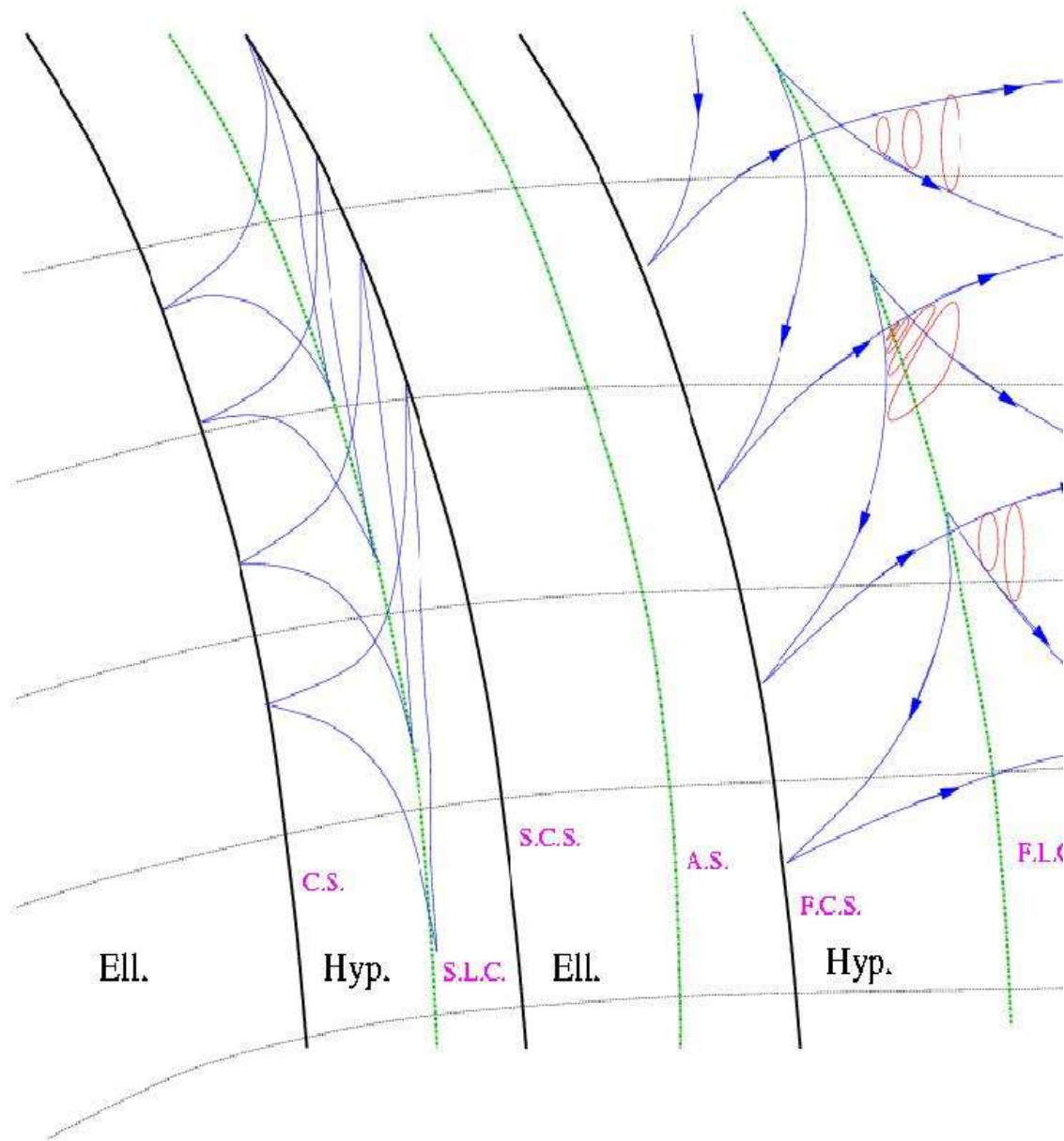
Hyp.

Axisymmetric Black Hole



Axisymmetric Wind





Self-similarity

- Separation of the variables: $A(r, \theta) \propto f(r) g(\theta)$
- Scaling of one of the variables with one coordinate

Partial → **Ordinary** diff. eqs.:

$$a_1 g(\theta) \frac{d^2 f(r)}{dr^2} + \frac{2a_2}{r} \frac{df(r)}{dr} \frac{dg(\theta)}{d\theta} + \frac{a_3}{r^2 \sin^2 \theta} f(r) \frac{d^2 g(\theta)}{d\theta^2} + a_4 + \frac{a_5}{1 - M^2} = 0$$

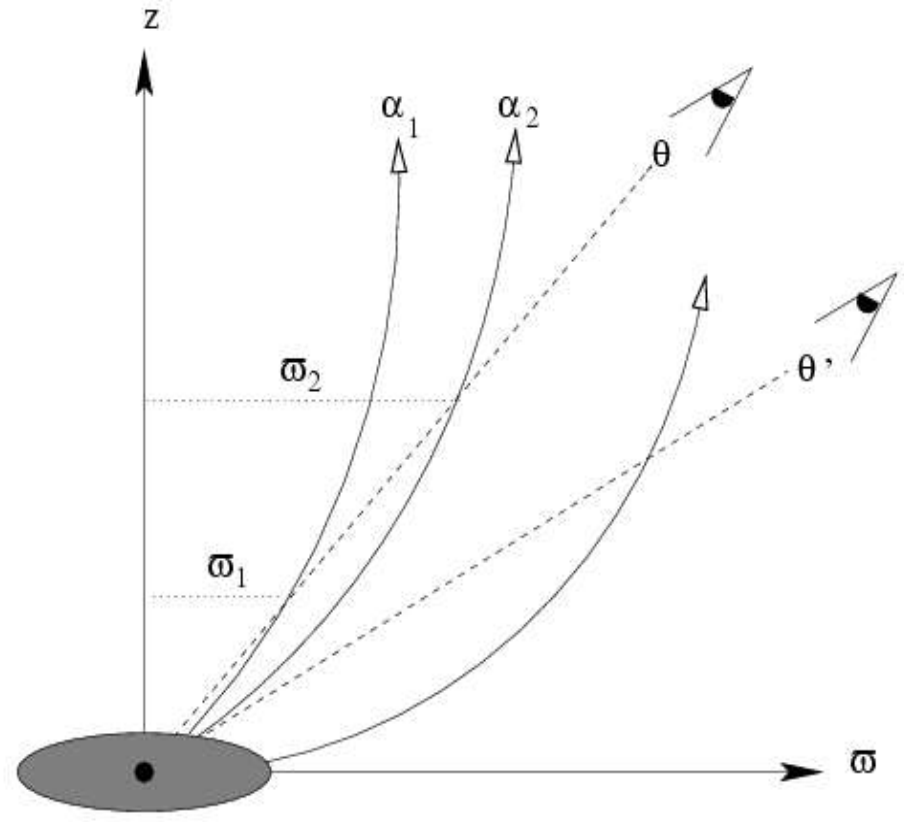
Radial: $f(r) \propto r^x$, $M(\theta)$

Singularities (L.C.) for $a_3 = 0$, where $V_\theta = V_{s,\theta}, V_{f,\theta}$
Magneto - Centrifugal Disk winds

Merid.: $g(\theta) \propto \sin^2 \theta$, $M(r)$

Singularities (L.C.) for $a_1 = 0$, where $V_r = V_{s,r}, V_{f,r}$
Magneto - Thermal winds

Radial self-similarity

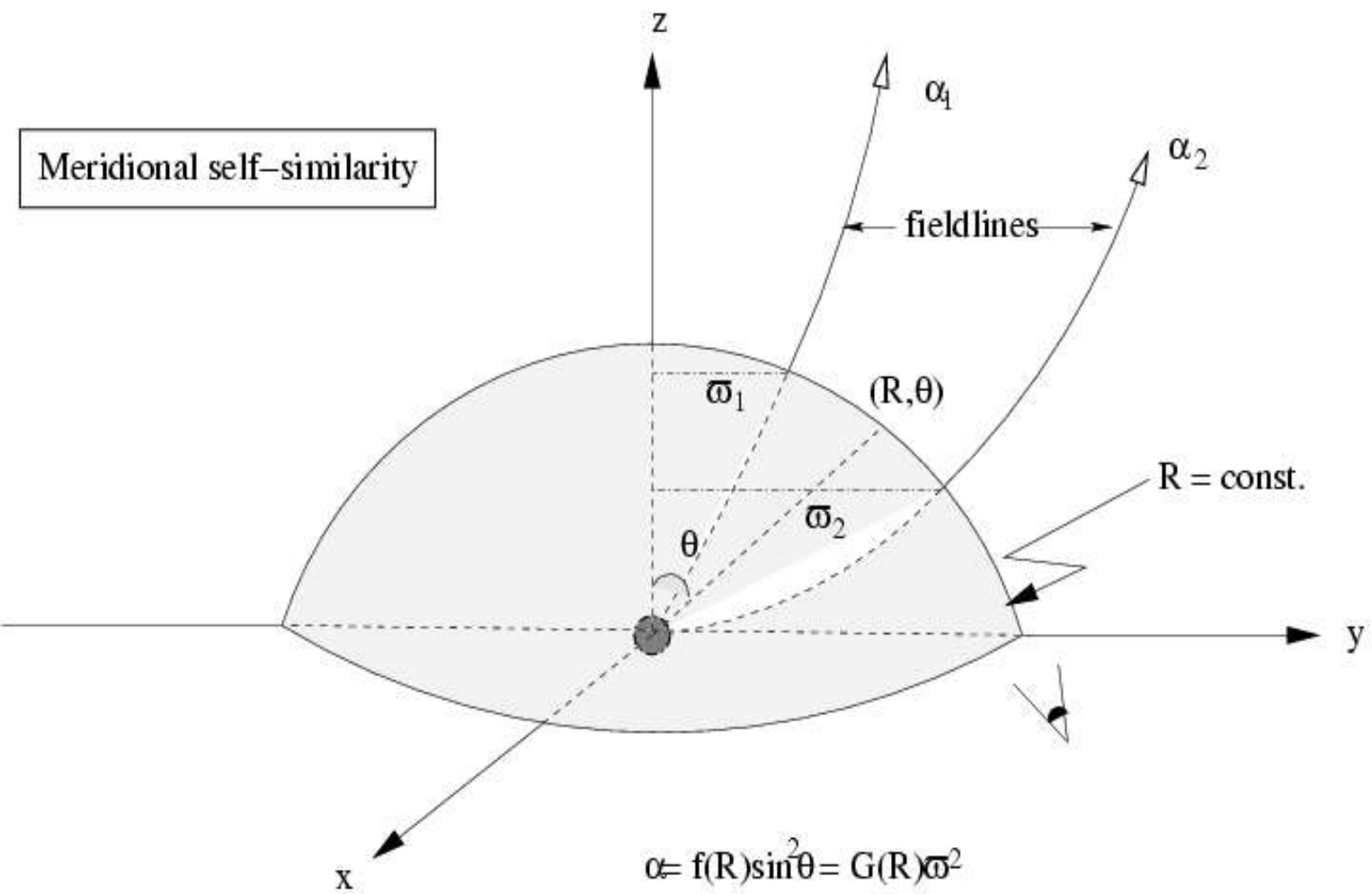


$$\alpha = g(\theta)R^x = F(\theta)\sigma^x$$

$$\theta = \text{const.} \longrightarrow \frac{\sigma_1}{\sigma_2} = \left(\frac{\alpha_1}{\alpha_2}\right)^{1/x}$$

Blandford & Payne '82,

Meridional self-similarity

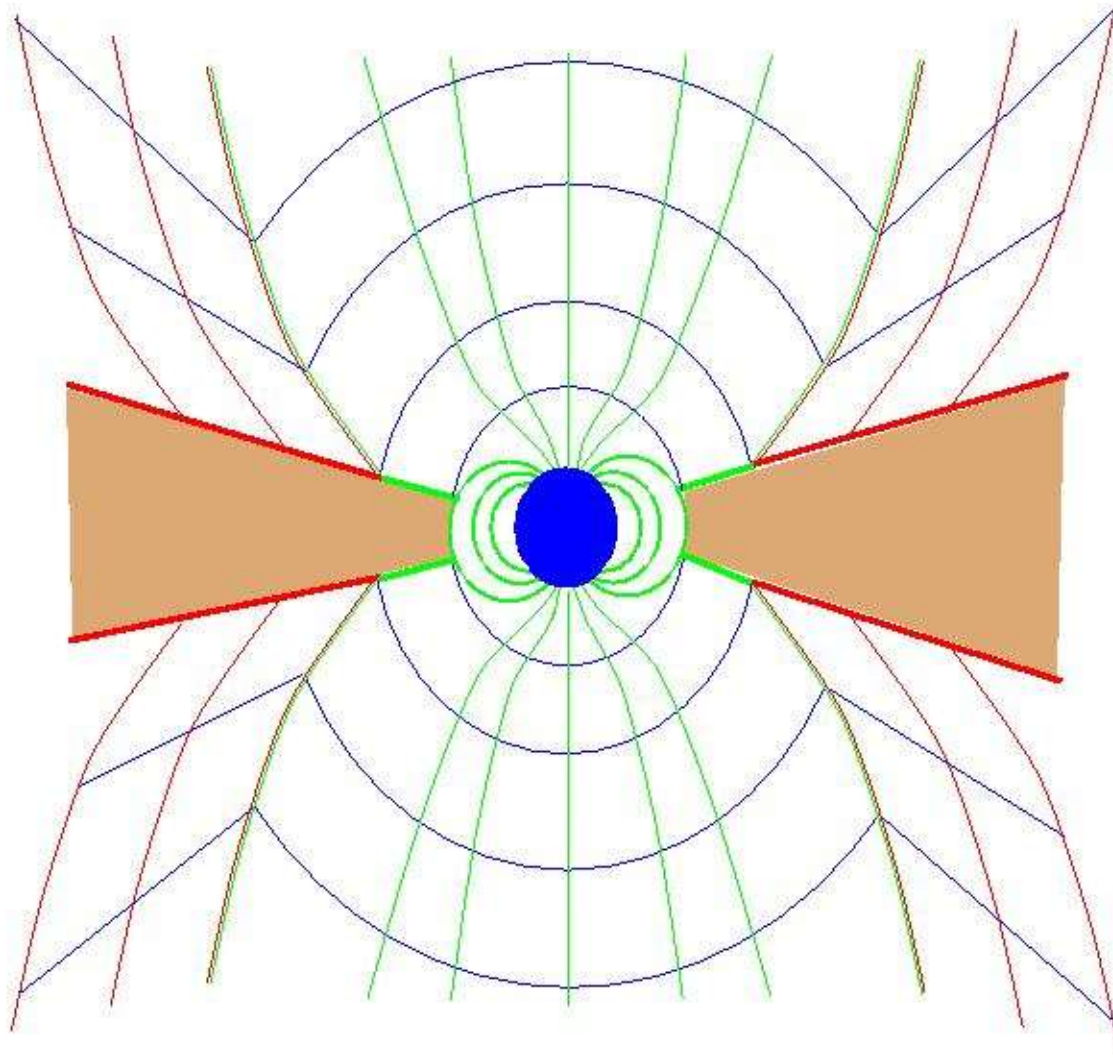


$$\alpha \propto f(R) \sin^2 \theta = G(R) \varpi^2$$

$$R = \text{const.} \longrightarrow \sqrt{\frac{\alpha_1}{\alpha_2}} = \frac{\varpi_1}{\varpi_2}$$

Low & Tsinganos '85, Sauty et al. '94,

SELF-SIMILAR JETS



Meridional self-similar models

(Meliani, Sauty, Tsinganos, Trussoni, Vlahakis)

$M(R) \rightarrow$ Surfaces of constant M : spheres

$$I_z \propto L\Psi_A \propto \lambda\alpha$$

$$P(R, \alpha) \propto P_o(R)(1 + \kappa\alpha)$$

$$\Psi_A^2(\alpha) \propto 1 + \delta\alpha$$

From numerical integration we have $\mathcal{F}(R; \lambda, \kappa, \delta, \dots)$

The solutions are suitable to model outflows in the region around **the rotational axis**

Acceleration: basically **thermal**

Collimation: **thermal** or **magnetic**

The **collimation** depends on the interplay between the inwards and outwards forces across the streamlines

The **thermal** or **magnetic** confinement is ruled by the following parameter

$$\epsilon \propto E_P - E_c + \Delta E_G$$

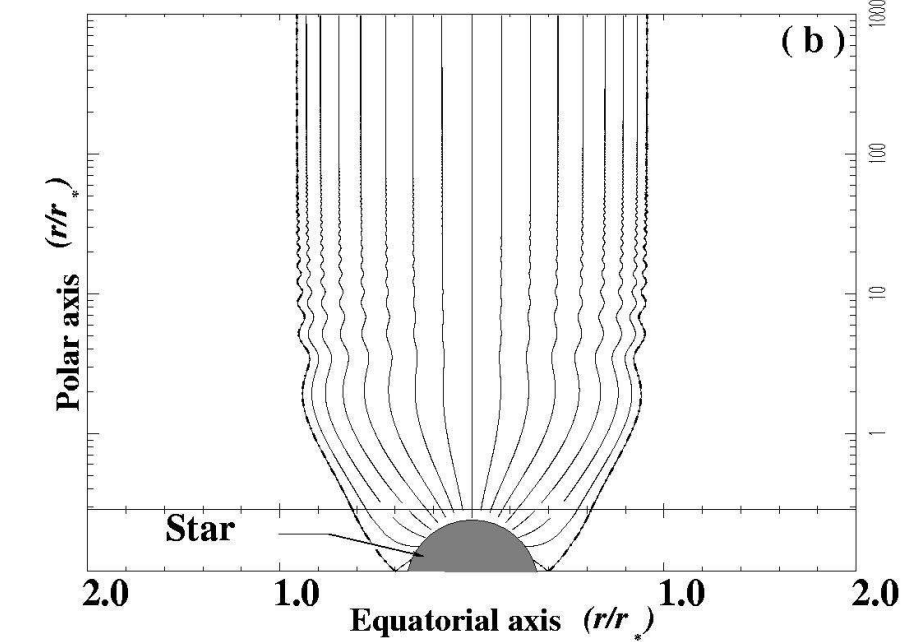
with

E_P : Variation of the Poynting energy

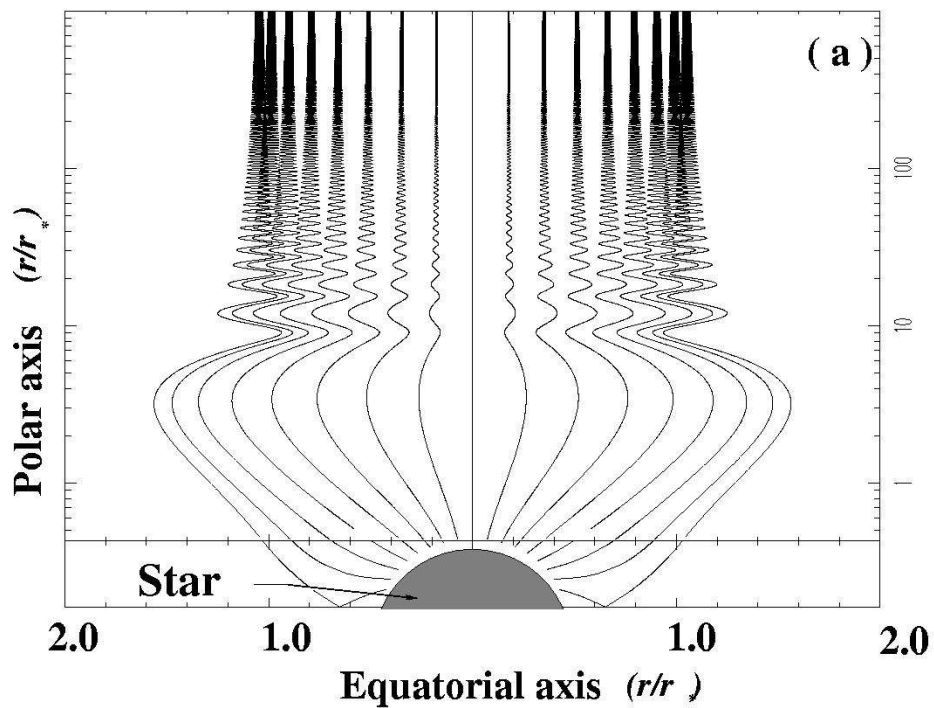
E_c : Variation of the centrifugal energy

$\Delta E_G \propto (\kappa - \delta)$: Variation of the gravitational energy

- $\epsilon > 0$ **Efficient** magnetic rotator: magn. collimation
- $\epsilon < 0$ **Inefficient** magn. rotator: thermal. coll.



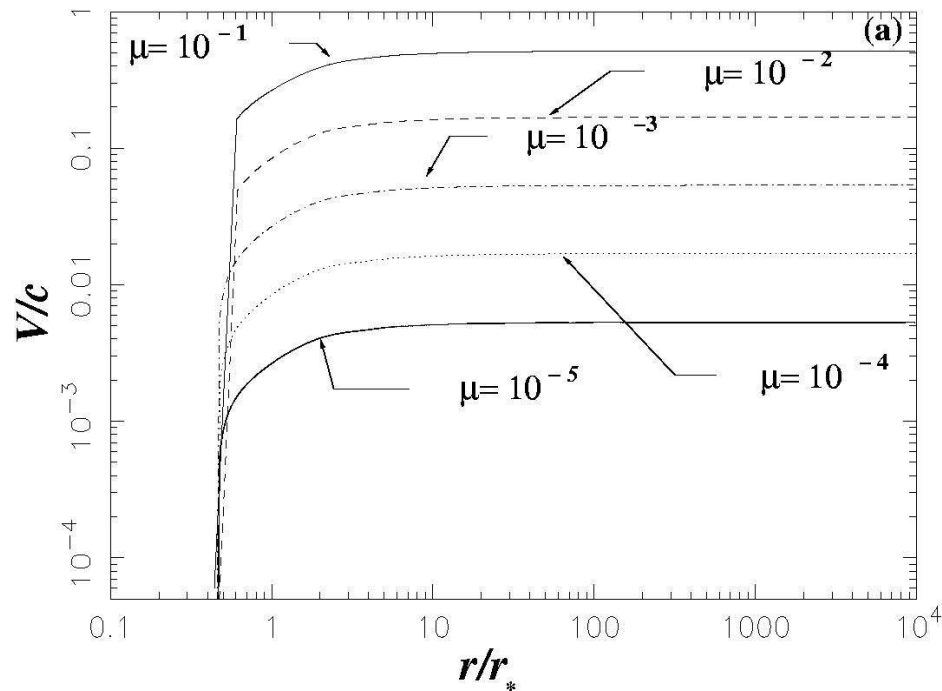
$$\varepsilon > 0$$

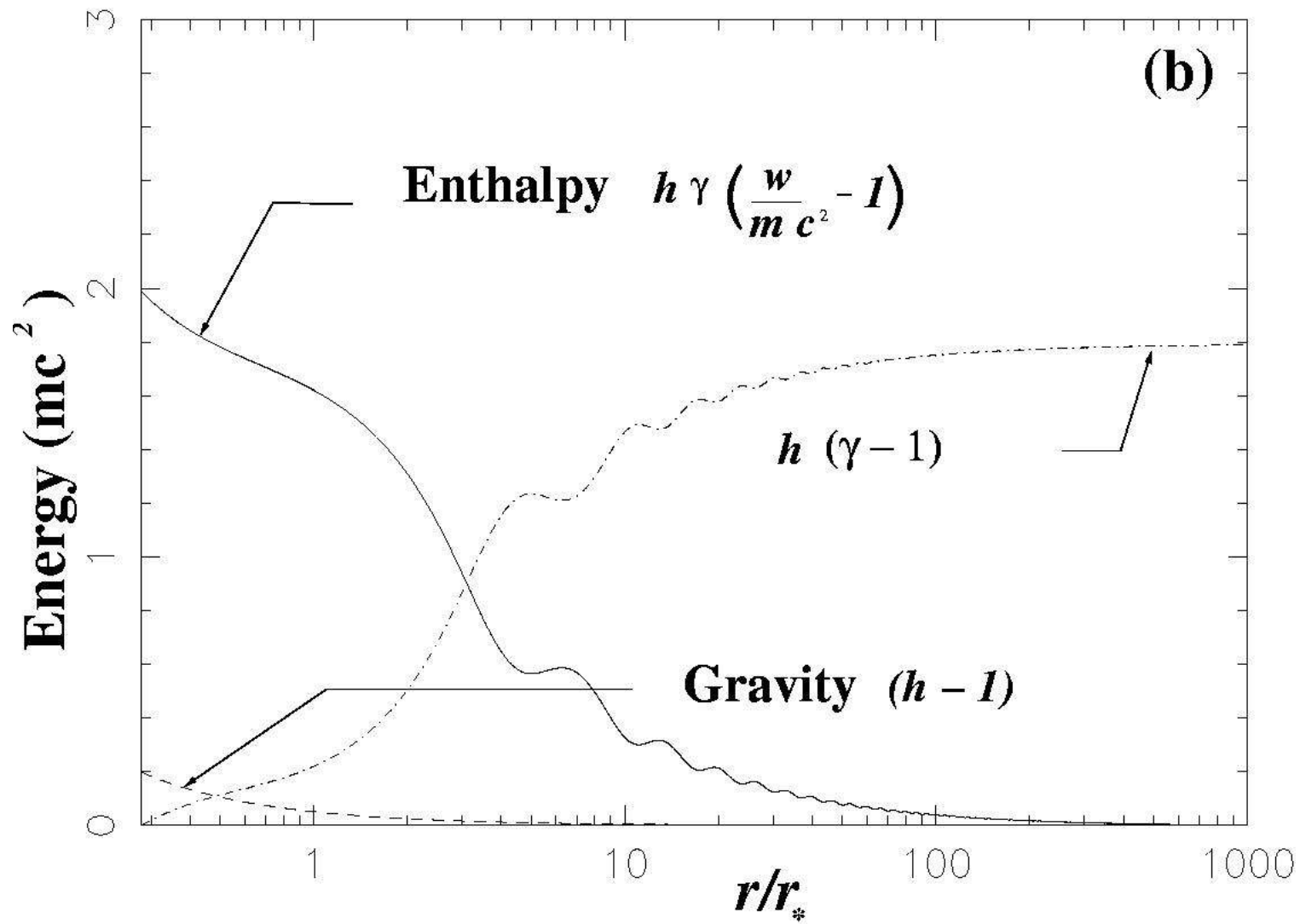


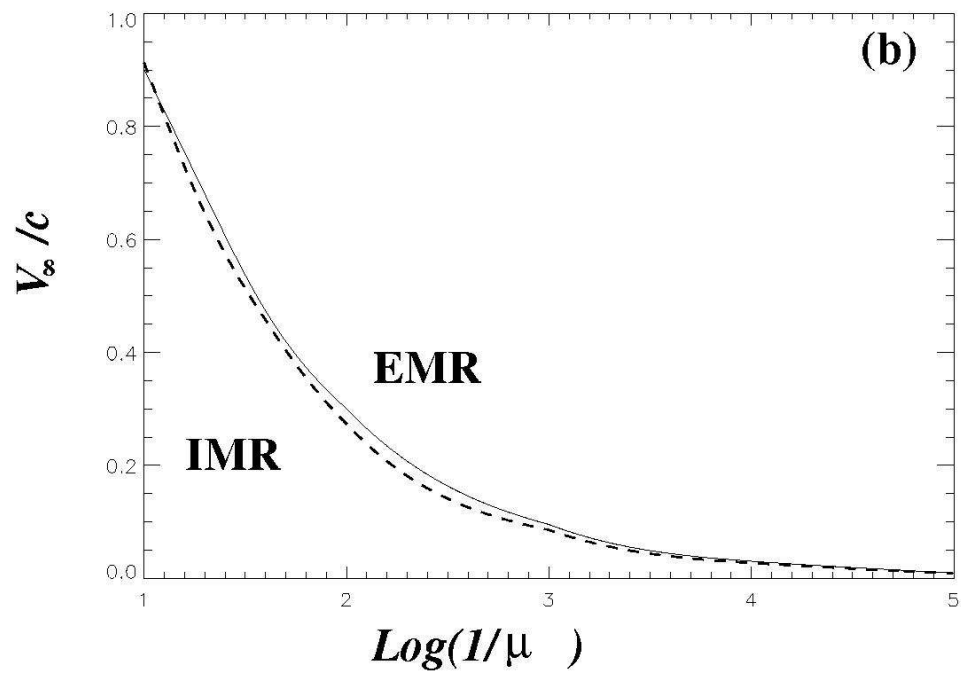
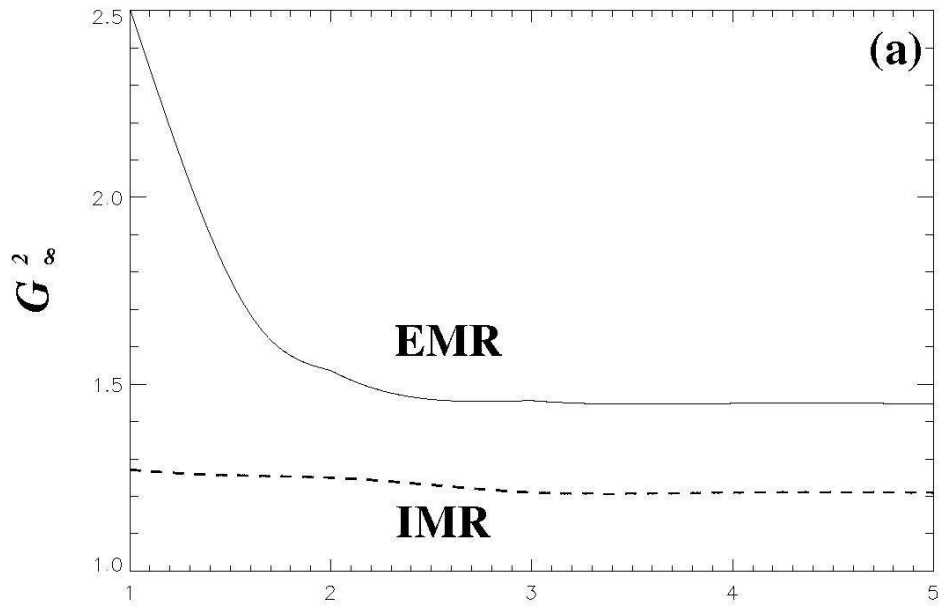
$$\varepsilon < 0$$

Relativistic Extension

- New parameter: $\mu = r_g/r_* < 1$
- Non relativistic rotational velocity: $\varpi \ll c/\Omega$
- *Mildly relat. velocities*
- *Electric force* $[\rho \mathbf{E} \propto -\rho (\mathbf{V} \times \mathbf{B}) / c]$: *decoll.*





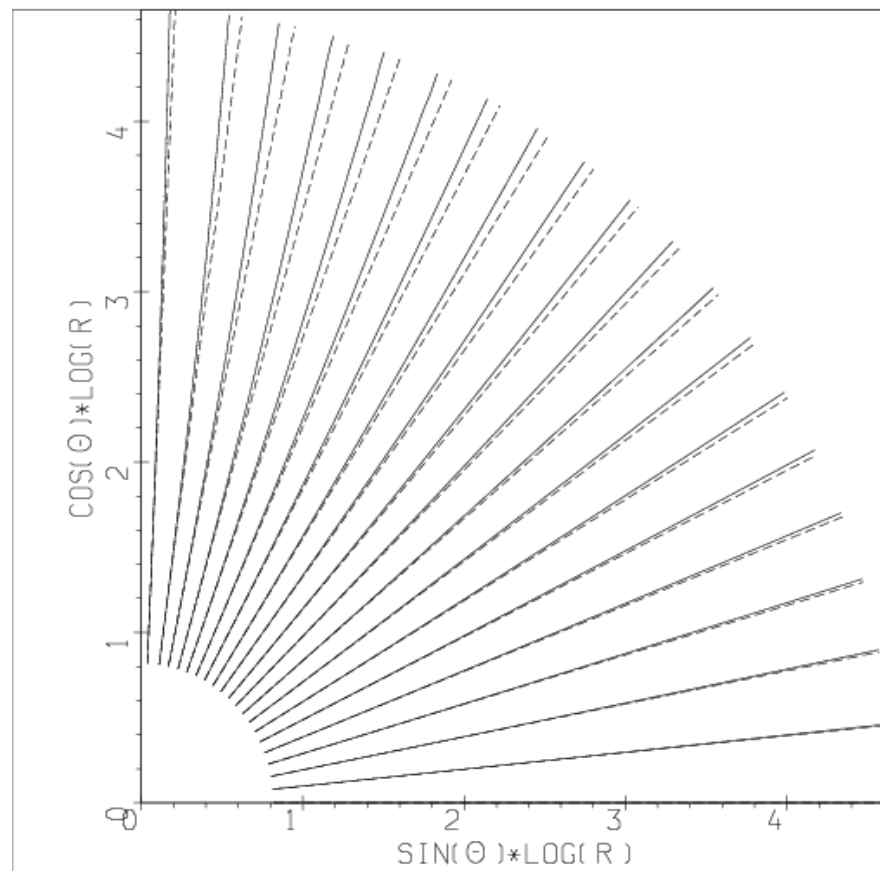


Weak collim. in relat. outflows
time dependent computations
(Tsinganos & Bogovalov,
2002-2005)

Possible Model:

Inner relativistic wind

Outer classical disk wind



CONCLUSIONS

Through the analysis of MHD steady solutions we can have insights on the physical properties of astrophysical outflows

Further issues:

Non ideal MHD (resistivity)

Radiative effects (losses and/or acceleration)

Large amplitude e.m. waves

Time dependent solutions: Asympt. status, instabilities
Interaction with the envir.
Launch

.....