# The jet-torus structure of Pulsar Wind Nebulae: relativistic MHD simulations

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### Outline

### • Pulsar Wind Nebulae in Supernova Remnants

- Observations
- Models (analytical, numerical)
- PWN inner jet-torus structure
  - Observations
  - Theoretical background
- PWN/SNR 2-D axisymmetric RMHD simulations
  - Overall dynamics, jet formation
  - Synchrotron emission and comparison with observations
- Summary and conclusions

### Papers on PWNe by our team

- Jet-torus in PWNe: synchrotron and polarization maps
  - Del Zanna, Volpi, Amato, Bucciantini, *in preparation*
  - Bucciantini, Del Zanna, Amato, Volpi, 2005, A&A, *submitted*
- Bow-shock PWNe
  - Bucciantini, Amato, Del Zanna, 2005, A&A, 434, 209
- Rayleigh-Taylor instabilities (filaments)
  - Bucciantini, Amato, Bandiera, Blondin, Del Zanna, 2004, A&A, 423, 253
- 2-D PWN-SNR simulations: jet-torus structure
  - Del Zanna, Amato, Bucciantini, 2004, A&A, 421, 1063
- 1-D PWN-SNR simulations
  - Bucciantini, Bandiera, Blondin, Amato, Del Zanna, A&A, 2004, 422, 609
  - Bucciantini, Blondin, Del Zanna, Amato, 2003, A&A, 405, 617
- RHD and RMHD numerical code
  - Del Zanna, Bucciantini, Londrillo, 2003, A&A, 400, 397
  - Del Zanna, Bucciantini, 2002, A&A, 390, 1177

# **Pulsar Wind Nebulae**



SNR

- PWNe (plerions) are hot bubbles emitting non-thermal radiation (synchrotron) at all wavelengths: require injection of relativistic particles and magnetic fields
- Originated by the interaction of the ultra-relativistic magnetized pulsar wind with the expanding SNR dense ejecta
- Crab Nebula in optical: central amorphous mass (continuum) + external filaments (lines)

# Sketch of PWN / SNR interaction



### Pulsar magnetosphere and wind



Coroniti, 1990

- Pulsar spin-down energy is converted to Poynting flux (mainly a toroidal field) and in a pair wind (with σ>>1)
- At the TS models predict σ<<1 to match the observed synchrotron emission: the sigma paradox!
- Striped wind: the magnetic field may decrease because of equatorial reconnection or dissipation of fast waves at TS

# PWN analytical MHD theory (KC84)

- PWN theory was mainly based on 1-D analytic (*Rees* & Gunn 1974; Kennel & Coroniti, 1984) and self-similar (*Emmering & Chevalier, 1987*) MHD models
- KC84 (spherically symmetric, stationary):
  - assume that the wind terminates with a strong MHD shock
  - solve the relativistic jump conditions at TS
  - solve the equations in the PWN region
  - calculate the synchrotron emission
  - a best fit analysis provides the wind parameters:

$$R_{TS} = 3 \times 10^{17} cm$$
,  $L = 5 \times 10^{38} erg/s$ ,  $\gamma = 3 \times 10^{6}$ ,  $\sigma = 3 \times 10^{-3}$ 

### Jet-torus structure: Chandra X-ray images



- Crab nebula (Weisskopf et al., 2000; Hester et al., 2002)
- Vela pulsar (Helfand et al., 2001; Pavlov et al., 2003)
- Other objects: PSR 1509-58, G0.9+01, G54.1+0.3

### Jet-torus structure: relativistic motions





Crab

Vela

- Equatorial motions (wisps): v=0.3-0.5 c
- Polar jet motions: v=0.5-0.8 c

### Jet-torus structure: theory

- Torus: higher equatorial energy flux
- Jets: magnetic collimation. But in PW:

$$\gamma > 1 \Longrightarrow \rho_q E + j \times B \approx 0$$

collimation downstream of the TS?

- Bogovalov & Khangoulian, 2002
- <u>Lyubarsky, 2002</u>
- Axisymmetric RMHD simulations of the interaction of an anisotropic relativistic magnetized wind with SN ejecta
  - Komissarov & Lyubarsky, 2003, 2004
  - Del Zanna, Amato & Bucciantini, 2004



# Axisymmetric relativistic wind model

- Far from the pulsar light cylinder the wind is expected to be ultrarelativistic, cold, and weakly magnetized. We assume:
  - Isotropic mass flux, anisotropic energy flux  $(F \propto r^2 \rho \gamma^2 \propto \gamma)$ :

 $\gamma(\theta) = \gamma_0 [\alpha + (1 - \alpha) \sin^2(\theta)]$ 

• Purely toroidal magnetic field (split monopole, Michel, 1973):

 $B(r,\theta) = B_0(r_0 / r) \sin(\theta)$ 

• Parameters of the wind model:

$$\gamma_{0} > 1, \alpha = \frac{F(0)}{F(\pi/2)} < 1, \sigma = \frac{B_{0}^{2}}{4\pi c^{2} \rho_{0} \gamma_{0}^{2}} < 1$$

# Simulation setup

- Central-type conservative RMHD code (HLL, second order)
- Spherical geometry, axial symmetry  $(r, \theta)$
- Poloidal velocity and purely toroidal magnetic field
- Computational grid: 400 points in r, 100 in  $\theta$
- Boundaries: injection for r=0.05 ly, extrapolation for r=20 ly
- Long time simulations (beginning of reverberation phase)
- High accuracy near the center: extremely small timesteps!
- Initial conditions:
  - Pulsar ultrarelativistic wind
  - Spherical shell of expanding dense ejecta
  - Static unmagnetized ISM

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### PWN self-similar evolution and TS shape



• Expected TS profile:

$$R(\theta) \approx R_0 \sqrt{\alpha} + (1-\alpha) \sin^2(\theta)$$



### **PWN elongation**

• Magnetic pinching effect (Begelman & Li, 1992):





### TS structure and flow pattern

- The wind anisotropy shapes the TS structure. A complex flow pattern arises:
  - A: ultrarelativistic pulsar wind
  - B: subsonic equatorial outflow
  - C: supersonic equatorial funnel
  - D: super-fastmagnetosonic flow
  - a: termination shock front
  - b: rim shock
  - c: fastmagnetosonic surface



### Formation of polar jets by hoop stresses

- The flow pattern changes drastically with increasing  $\sigma$
- For high magnetization ( $\sigma$ >0.01) a supersonic jet is formed



### Dependence on the field shape

Initial magnetic field with a narrow equatorial neutral sheet



# A model for synchrotron emission

- How to build synchrotron emission maps:
  - Assume a power law spectrum of electron energies at TS

$$f_{0}(\boldsymbol{\varepsilon}_{0}) \propto p_{0} \boldsymbol{\varepsilon}_{0}^{-(2\alpha+1)}$$

• Evolve the energy considering adiabatic and synchrotron losses

$$\frac{d\varepsilon}{dt'} = \varepsilon \frac{d}{dt'} \ln(\rho^{1/3}) - \frac{4e^4}{9m^3c^5} (B')^2 \varepsilon^2$$

• Assume emission at the critical frequency

$$v \propto B_{\perp}' \varepsilon^2$$

• Calculate the spectral emissivity function in the observer frame

$$j_{\nu}(\mathbf{v},\hat{n}) \propto \gamma D^{3+\alpha} p(B'_{\perp})^{\alpha+1} [1-\varepsilon(\mathbf{v})/\varepsilon_{\infty}]^{2\alpha-1} \mathbf{v}^{-\alpha}$$

• Obtain synthethic maps by integrating along the LOS

### Comparison with observations: maps

- Effects of synchrotron losses: optical vs X-ray maps
  - Runs with expanding CD at given velocity and realistic luminosity X-ray



### Comparison with observations: maps

- Constraining the field shape of the pulsar wind:
  - Runs with narrower *striped wind* region reproduce observations better



σ=0.03. b=1

### σ=0.03. b=10

#### L. Del Zanna: The jet-torus structure of PWNe

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### Comparison with observations: maps

Simulated X-ray maps vs Chandra images:



### Comparison with observations: spectrum

Synchrotron spectral index X-ray maps:







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### Comparison with observations: polarization

- Simulated optical high resolution polarization maps
  - A toy model first: uniform emitting torus



## Comparison with observations: polarization

- Simulated optical high resolution polarization maps
  - Results from the relativistic MHD simulations



## Summary and conclusions

- Many PWNe show a jet-torus structure (Crab, Vela, ...)
- The torus is explained with a higher equatorial energy flux
- Jet collimation forbidden in the wind. Inside PWN?
- RMHD axisymmetric simulations confirm this scenario:
  - The TS has a toroidal shape, a strong equatorial flow is produced
  - For  $\sigma$ >0.01 hoop stresses divert the flow toward the axis
  - Plasma is compressed and a polar jet with v=0.5-0.7c is launched
  - Simulated synchrotron maps resemble closely X-ray images
  - Work in progress: constraining B, spectra and polarization maps

### Thank you